

February 4, 2004 Lecture 1

18.306 - Advanced PDE with Applications

Dionisios Margetis

ODE

Ex.1 ODE: $\frac{du}{dx} = 0, u = u(x)$

solution is $u = C$, all x

Condition: $u(x_0) = u_0 \Rightarrow C = u_0$

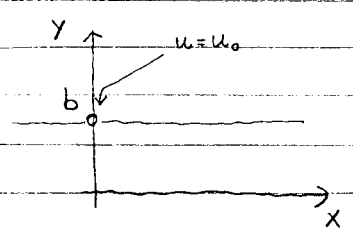
$\Rightarrow u = u_0$ for all x

Conclusion ODE's involve arbitrary constants.

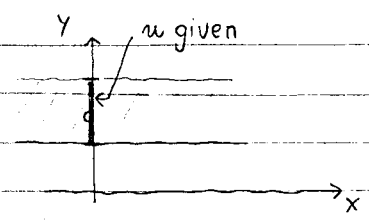
Ex.2 PDE: $\frac{\partial u}{\partial x} = 0, u = u(x, y)$

solution is $u(x, y) = C(y)$ "arbitrary function of y "

General solutions to PDE's involve arbitrary functions.

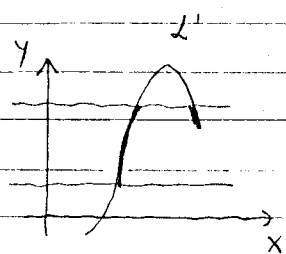
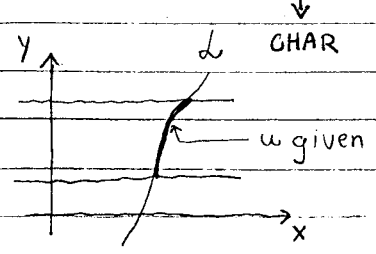


- by knowing the solution at $(0, b)$ we know the solution at the horizontal line passing through $(0, b)$



- by knowing u on the segment $y_1 < y < y_2$ we know the solution in the whole strip $y_1 < y < y_2$, all x

$y = \text{const}$: "CHARACTERISTIC" by def. line on which u is constant



not possible if u has to be continuous

If u given on ∂' then for some $y, y_1 < y < y_2$ I have more than one value of w ,

→ not OK unless I allow for "jumps" in u

★ More, generally, along a CHAR $PDE \Rightarrow ODE$ THEME FOR THE COURSE

★ If a PDE reduces to ODE, it is considered solvable.

Initial Value Problem (IVP) if values of w are given on some open, smooth curve (eg. ∂')

Generally, PDE is an equ of the form

$$F(x, y, w, u_x, u_y, u_{xx}, u_{xy}, u_{yy}, \dots) = 0 \quad \text{where } u = u(x, y)$$

Order of the PDE is the order of the highest derivative of w .

Linear PDE: F is linear function of w and its derivatives.

Ex. 3 PDEs $c = \text{const}$

- $u_t - c u_x = 0$ (kinematic eqn)
traffic flow, gas dynamics
- $u_{tt} - c^2 u_{xx} = 0$ (wave eqn)
- $u_{xx} + u_{yy} = 0$ (Laplace eqn)
 $\nabla^2 u = 0$

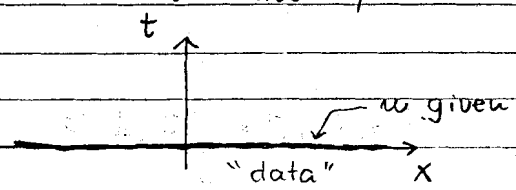
The issue of conditions for PDE

Ex. 4 $\begin{cases} u'' + u = 0, & u = u(x) \quad 0 \leq x \leq 1 \\ u(0) = 0 \\ u'(0) = 1 \end{cases}$ IVP for ODE

Ex. 5 $\begin{cases} u'' + u = 0, & u = u(x) \quad 0 \leq x \leq 1 \\ u(0) = 0 \\ u(1) = 1 \end{cases}$ BVP for ODE

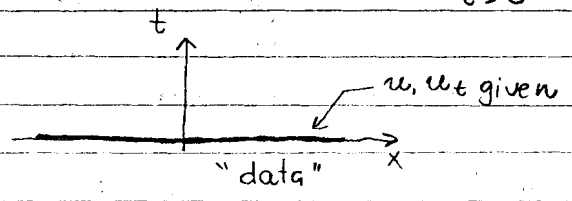
Ex. 6 $u_t - cu_x = 0, \quad -\infty < x < +\infty$
 $t > 0$

conditions are usually



IVP (there is a unique solution)

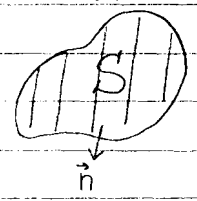
Ex. 7 $u_{tt} - c^2 u_{xx} = 0, \quad -\infty < x < +\infty$
 $t > 0$



IVP (there is a unique solution)

"Cauchy data" = u and its normal derivative (u_t) are given on a line

Ex. 8 $\nabla^2 u = 0 \quad \partial S$



- (A) u given on ∂S BVP "Dirichlet"
- (B) $\frac{\partial u}{\partial n}$ given on ∂S BVP "Neumann"
- (C) $au + b\frac{\partial u}{\partial n}$ given on ∂S BVP "Robin" (mixed bc's)

→ for these conditions (ex. 6, 7, 8) solutions not only exist but are unique.

In physical problems, solutions: exist
be unique }

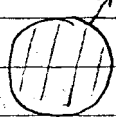
Ex. 9 $\begin{cases} u'' = 0 \\ u'(0) = 0, u'(1) = 1 \end{cases}, \quad 0 < x < 1 \quad u = u(x)$

$u'' = 0 \rightarrow u(x) = Ax + B, \quad u'(x) = A$

- $u'(0) = 0 \Rightarrow A = 0$
 - $u'(1) = 1 \Rightarrow A = 1$
- must to be true together \Rightarrow no solution
 Impossible

Alternatively: $0 = \int_0^1 dx u''(x) = u'(1) - u'(0) = 1 - 0 = 1$! contradiction

Ex. 10 $\nabla^2 u = 0$

BVP:  $\frac{\partial u}{\partial r} = 1$ Neumann bc's

Recall: Divergence theorem $\int_S \nabla \cdot \vec{A} d\vec{S} = \int_C \hat{n} \cdot \vec{A} dl$

$0 = \int_S \nabla \cdot \nabla u d\vec{S} = \int_C \hat{n} \cdot \nabla u dl = \int_0^{2\pi} \frac{\partial u}{\partial r} dl = 1 \times 2\pi r$
contradiction!
 \Rightarrow no solution with this data

ill-posed problems = no solution exists for the given eqn with the given data

Ex. 11 $\nabla^2 u = 0$ BVP: $\frac{\partial u}{\partial r} = 0$ no contradiction this time

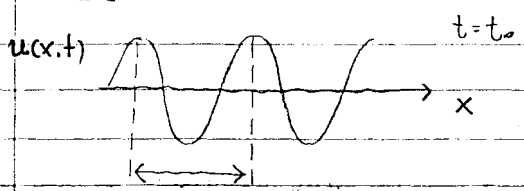
let u_1 is a solution then $u_2 = u_1 + \text{const}$ is also a solution
 \Rightarrow solution exists but it is not unique

Well-posed problems (classical definition by Hadamard):

- 1) the solution must exist
- 2) the solution must be unique
- 3) the solution depends continuously on data
(eg. small perturbation of the data results in small perturbation of soln)

Concepts that apply both to linear and nonlinear PDE's

Wave



λ wavelength

$k = \frac{2\pi}{\lambda}$ wave number (how many crests/length)

$u(x,t) = A \cos(\omega t - kx)$, ω frequency (in time)

satisfies the wave equation $u_{tt} - c^2 u_{xx} = 0$, $c = \frac{\omega}{k} = \omega \lambda$

