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# Engineering Analysis with Boundary Elements

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## Bipolar coordinates, image method and the method of fundamental solutions for Green's functions of Laplace problems containing circular boundaries

J.T. Chen<sup>a,b,\*</sup>, H.C. Shieh<sup>a</sup>, Y.T. Lee<sup>a</sup>, J.W. Lee<sup>a</sup><sup>a</sup> Department of Harbor and River Engineering, National Taiwan Ocean University, Keelung 20224, Taiwan<sup>b</sup> Department of Mechanical and Mechatronic Engineering, National Taiwan Ocean University, Keelung 20224, Taiwan

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### ABSTRACT

Green's functions of Laplace problems containing circular boundaries are solved by using analytical and semi-analytical approaches. For the analytical solution, we derive the Green's function using the bipolar coordinates. Based on the semi-analytical approach of image method, it is interesting to find that the two frozen images for the eccentric annulus using the image method are located on the two **foci** in the bipolar coordinates. This finding also occurs for the cases of a half plane with a circular hole and an infinite plane with two circular holes. The image method can be seen as a special case of the method of fundamental **solutions**, which only at most four unknown strengths are required to be determined. The optimal locations of sources in the method of fundamental solutions can be captured using the image method and they are dependent on the source location and the geometry of problems. Three illustrative examples were demonstrated to verify this point. Results are satisfactory.

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### 1. Introduction

A number of physical and engineering problems governed by the Laplace equation in two independent variables, e.g., steady-state heat conduction, electrostatic potential and fluid flow, were solved using the conformal mapping to obtain an analytical solution. Besides, we can formulate the same problems using special curvilinear coordinates to obtain a solution, e.g., bipolar coordinates and elliptic coordinates. Carrier and Pearson [1] employed the bilinear transformation of conformal mapping to solve certain kinds of potential problems. An eccentric case was mapped to an annular domain through a bilinear transformation. For a polygonal shape, it can also be mapped to a regular region using the Schwarz–Christoffel transformation [2]. For a regular geometry, it is easy to solve the Laplace problem using the polar or Cartesian coordinates. Muskhelishvili [3] gave us a detailed description how an eccentric annulus can be mapped into concentric annulus using a simple form of linear fractional transformation. Chen and Weng [4] also used a similar method to solve eccentric annulus problems. Although a bilinear transformation was used, the mapping functions were not exactly the same between the one of Carrier and Pearson [1] and that of Muskhelishvili [3]. Problems of eccentric annulus, a half plane with a circular hole or an infinite plane containing two circular

holes usually use the bipolar coordinates to derive the analytical solution [5]. Ling [6], Timoshenko and Goordier [7], and Lebedev et al. [8] all presented an analytic solution by using the bipolar coordinates for the torsion of an eccentric bar. However, the mapping functions were not exactly the same. One is a cotangent function [6], another is a hyperbolic tangent function [8] and the other is a hyperbolic cotangent function [7]. After the bipolar coordinate system is introduced, the problem of special domain can be solved by using the separation of variables. Although Carrier and Pearson [1], Muskhelishvili [3], Ling [6], Timoshenko and Goordier [7] have solved the eccentric Laplace problems, their approaches are very similar, but not identical. Chen et al. [9] found that all the above-mentioned approaches can be unified after suitable transformations, translation, rotation and taking *log* in the conformal mapping. However, we will focus on Green's function instead of BVP without sources [10] in this paper.

Green's function has been studied and applied in science and engineering by mathematicians as well as engineers, respectively [11]. A computer-friendly solution for the potential generated by a point source in the ring-shaped region was studied by Melnikov and Arman [12]. In order to derive Green's function, Thomson [13] proposed the concept of reciprocal radii to find the image source to satisfy the homogeneous Dirichlet boundary condition using the image method. Greenberg [14] and Riley et al. [15] employed a trick to satisfy the boundary condition for two special points, then the image location can be determined. Chen and Wu [16] proposed a natural and logical way to find the location of image and the strength by employing the degenerate kernel. The image method is a classical approach for constructing Green's function. In certain cases, it is possible to obtain the exact solution for

\* Corresponding author at: Department of Harbor and River Engineering, National Taiwan Ocean University, Keelung 20224, Taiwan.

Tel.: +886 2 24622192x6177; fax: +886 2 24632375.

E-mail address: jtchen@mail.ntou.edu.tw (J.T. Chen).

a concentrated source in a bounded domain through superimposing the infinite plane solution for the given source and its image source. Although the scope of this method is limited, it yields a great deal of insight into the solution when it works [17]. Here, we will extend to a semi-analytical approach once the closed-form solution using the image method is not possible. Our goal is to broaden the scope of applications on the image method.

In this paper, we have three issues. First, the image method is seen as a special case of the method of fundamental solutions since its image singularities locate outside the domain. Second, the optimal locations of the method of fundamental solutions sources are found to be dependent on the source location and the geometry of the problems. Third, it is found that the two frozen images of the image method are located on the two focuses in the bipolar coordinates. Using the bipolar coordinates and the image method, three cases, an eccentric annulus, a half plane with a circular hole and an infinite domain with two holes, are solved. The bipolar coordinates are reviewed for the eccentric ring in Section 2. In Section 3, the image method is employed to derive Green's function for problems containing circular boundaries. Numerical results are given in Section 4. Finally, a conclusion is drawn in Section 5.

**2. Geometric characterization of the bipolar coordinates**

The relation between the bipolar coordinates  $(\xi, \eta)$  and the Cartesian coordinates  $(x, y)$  [9] is defined by

$$x + iy = icc \cot\left(\frac{1}{2}\zeta\right), \quad \zeta = \xi + i\eta, \tag{1}$$

where  $c$  is a positive constant. Eq. (1) yields

$$x = c \frac{\sinh \eta}{\cosh \eta - \cos \xi}, \quad y = c \frac{\sin \xi}{\cosh \eta - \cos \xi}, \tag{2}$$

where  $-\pi \leq \xi < \pi, -\infty < \eta < \infty$ . By eliminating  $\xi$  in Eq. (2), we obtain a circle with the center at  $(c \coth \eta, 0)$  and the radius  $ccsch \eta$  as follows:

$$(x - c \coth \eta)^2 + y^2 = c^2 csch^2 \eta. \tag{3}$$

Elimination of  $\eta$  from Eq. (2) results in the other circle with the center at  $(0, c \cot \xi)$  and the radius of  $ccsc \xi$  as follows:

$$x^2 + (y - c \cot \xi)^2 = c^2 csc^2 \xi. \tag{4}$$

**Q3** Denoting by  $(\Gamma_1, \phi_1)$  and  $(\Gamma_2, \phi_2)$ , we have

$$x + iy + c = \Gamma_1 e^{i\phi_1}, \quad x + iy - c = \Gamma_2 e^{i\phi_2}, \tag{5}$$

$$\eta = \log(\Gamma_1 / \Gamma_2), \quad \xi = \phi_2 - \phi_1. \tag{6}$$

It follows that a curve  $\xi = \text{constant}$  is a family of circles passing through the poles  $(\pm c, 0)$ . The curve of  $\eta = \text{constant}$  shows a curve for which  $\Gamma_1 / \Gamma_2 = \text{constant}$ . The eccentric annulus is shown in Fig. 1. The outer radius  $b$ , inner radius  $a$  and the distance  $d$  are determined from Eq. (3) as follows:

$$a = ccsch(\eta_1), \tag{7}$$

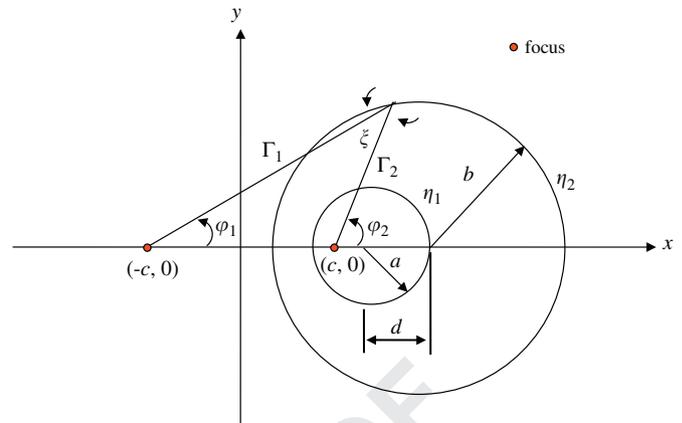


Fig. 1. Geometry relation of bipolar coordinates.

$$b = ccsch(\eta_2), \tag{8}$$

$$d = c[\coth(\eta_2) - \coth(\eta_1)]. \tag{9}$$

To describe an eccentric annulus in the bipolar coordinates, the three parameters,  $c, \eta_1$  and  $\eta_2$  are determined as follows:

$$c = \frac{\sqrt{a^4 + b^4 - 2a^2b^2 - 2d^2(a^2 + b^2) + d^4}}{2d}, \tag{10}$$

$$\eta_1 = \sinh^{-1}\left(\frac{c}{a}\right), \tag{11}$$

$$\eta_2 = \sinh^{-1}\left(\frac{c}{b}\right), \tag{12}$$

where  $\eta_1$  and  $\eta_2$  denote the inner and outer circles, respectively. Then, we can describe an eccentric annulus using the bipolar coordinates. In this case, Green's function was derived in terms of the bipolar coordinates as shown below [5]

$$G(\xi, \eta; \xi_0, \eta_0) = \begin{cases} \frac{1}{2\pi} \left[ \frac{(\eta_1 - \eta)(\eta_2 - \eta_0)}{\eta_1 - \eta_2} + 2 \sum_{n=1}^{\infty} \frac{\sinh n(\eta_1 - \eta) \sinh n(\eta_2 - \eta_0)}{n \sinh n(\eta_1 - \eta_2)} \cos n(\xi - \xi_0) \right], & \eta_1 \geq \eta \geq \eta_0, \\ \frac{1}{2\pi} \left[ \frac{(\eta - \eta_2)(\eta_0 - \eta_1)}{\eta_1 - \eta_2} + 2 \sum_{n=1}^{\infty} \frac{\sinh n(\eta_2 - \eta) \sinh n(\eta_1 - \eta_0)}{n \sinh n(\eta_1 - \eta_2)} \cos n(\xi - \xi_0) \right], & \eta_0 \geq \eta \geq \eta_2, \end{cases} \tag{13}$$

where  $(\xi_0, \eta_0)$  is the position of the source point.

**3. Image method**

For a problem of two-dimensional eccentric annulus as shown in Fig. 2, Green's function  $G(x, s)$  satisfies

$$\nabla^2 G(x, s) = \delta(x - s), \quad x \in \Omega, \tag{14}$$

where  $\Omega$  is the domain of interest,  $x$  the field point and  $\delta$  denotes the Dirac-delta function for the source at  $s$ . For simplicity, Green's function is considered to be subject to the homogeneous Dirichlet boundary conditions. In this case, we obtain the location of image point using the fundamental solution and matching the boundary condition. The eccentric annulus can be seen as a combination of interior and exterior problems as shown in Fig. 3. The source point and the image point are  $s$  and  $s'$  in Fig. 3, respectively. By matching the homogeneous Dirichlet boundary conditions for the interior or exterior boundaries the, position of the image source is at  $(a^2/R_s, \theta)$ , where  $s = (R_s, \theta)$ . We consider the fundamental solution

$U(x, s)$  for the infinite plane that is governed by

$$\nabla^2 U(x, s) = 2\pi\delta(x-s). \tag{15}$$

The closed-form fundamental solution is given as

$$U(x, s) = \ln r, \tag{16}$$

where  $r$  is the distance between  $s$  and  $x(r \equiv |x-s|)$ . Based on the separable property of the addition theorem or the so-called degenerate kernel, the fundamental solution  $U(x, s)$  can be expanded into a series form by separating the field point  $x(\rho, \phi)$  and source point  $s(R_s, \theta)$  in the polar coordinates

$$U(x, s) = \begin{cases} U^I(\rho, \phi; R_s, \theta) = \ln R_s - \sum_{m=1}^{\infty} \frac{1}{m} \left(\frac{\rho}{R_s}\right)^m \cos m(\theta - \phi), & R_s \geq \rho, \\ U^E(\rho, \phi; R_s, \theta) = \ln \rho - \sum_{m=1}^{\infty} \frac{1}{m} \left(\frac{R_s}{\rho}\right)^m \cos m(\theta - \phi), & R_s < \rho. \end{cases} \tag{17}$$

The image method can solve Green's function of eccentric case in a semi-analytical manner. Following the successive image process, it is found that the final two image locations freeze at the  $s_{c1}$  and  $s_{c2}$ . For the eccentric case, the distance from the center of outer circle to the source is  $R_s = \sqrt{x_0^2 + y_0^2}$ . The successive former

four locations of images are

$$\begin{aligned} R_1 &= \frac{b^2}{R_s}, & x_1 &= R_1 \cos \theta_1, & y_1 &= R_1 \sin \theta_1, \\ R_2 &= \frac{a^2}{\sqrt{(x_0+d)^2 + y_0^2}}, & x_2 &= R_2 \cos \theta_2 - d, & y_2 &= R_2 \sin \theta_2, \\ R_3 &= \frac{b^2}{\sqrt{x_2^2 + y_2^2}}, & x_3 &= R_3 \cos \theta_3, & y_3 &= R_3 \sin \theta_3, \\ R_4 &= \frac{a^2}{\sqrt{(x_1+d)^2 + y_1^2}}, & x_4 &= R_4 \cos \theta_4 - d, & y_4 &= R_4 \sin \theta_4, \end{aligned} \tag{18}$$

$$\begin{aligned} \theta_1 &= \theta, \\ \theta_2 &= \tan^{-1}\left(\frac{y_2}{x_2+d}\right), \\ \theta_3 &= \tan^{-1}\left(\frac{y_3}{x_3}\right), \\ \theta_4 &= \tan^{-1}\left(\frac{y_4}{x_4+d}\right), \end{aligned} \tag{19}$$

Green's functions for the two cases: (a) an eccentric annulus and (b) a half plane with a circular hole, can be represented by

$$G(x, s) = \frac{1}{2\pi} \left\{ \ln|x-s| - \lim_{N \rightarrow \infty} \left[ \sum_{i=1}^N (\ln|x-s_{4i-3}| + \ln|x-s_{4i-2}| - \ln|x-s_{4i-1}| - \ln|x-s_{4i}|) + c_1(N) \ln|x-s_{c1}| + c_2(N) \ln|x-s_{c2}| + e(N) \right] \right\}, \tag{20}$$

where  $s_{4i-3}$ ,  $s_{4i-2}$ ,  $s_{4i-1}$  and  $s_{4i}$  are the successive image locations [18],  $e(N)$  can be understood as a rigid body term,  $c_1(N)$  and  $c_2(N)$  are the singularity strengths of the two frozen points at  $s_{c1}$  and  $s_{c2}$ , which can be determined by matching the boundary conditions. For the case of infinite plane with two holes, the expression of Green's function is given in Table 1. Table 1 demonstrates that the frozen image points  $s_{c1}$  and  $s_{c2}$  happen to be the focuses in the bipolar coordinates.

4. Illustrative examples

Case 1. An eccentric case (a special case: annular case [12])

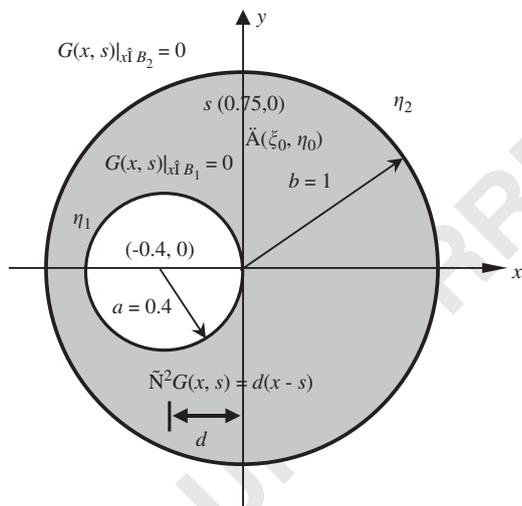


Fig. 2. Problem sketch for Green's function of an eccentric annulus.

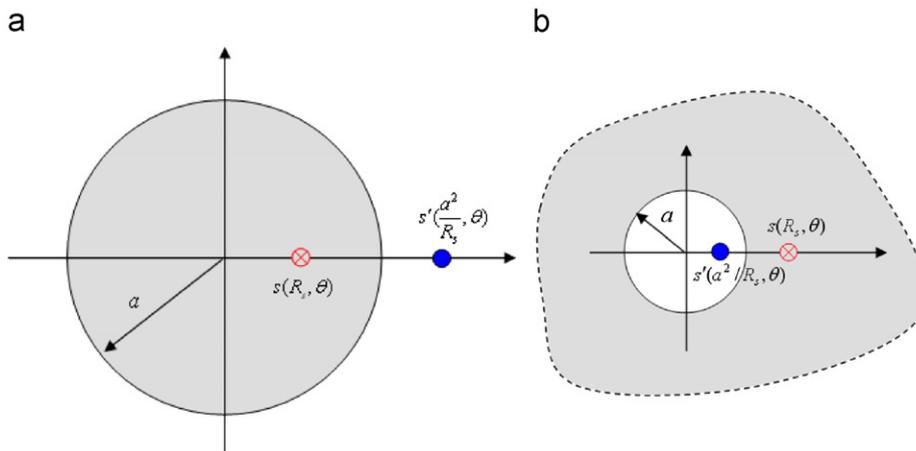
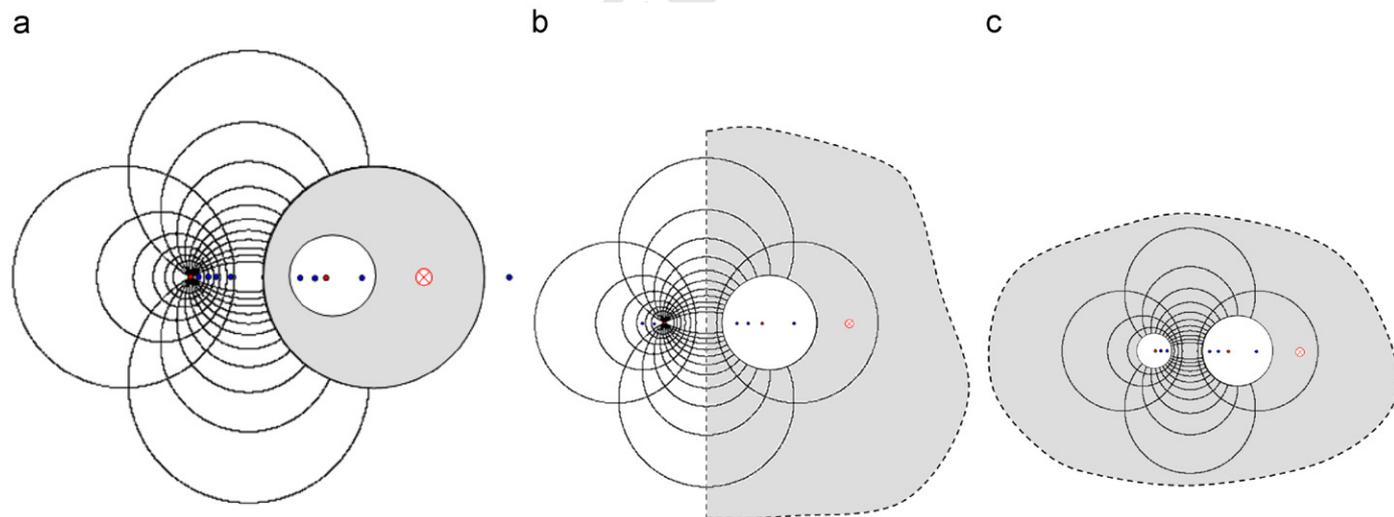


Fig. 3. Sketch of position of image point (a) an interior case and (b) an exterior case.

**Table 1**  
Frozen points of the image method and focuses in the bipolar coordinates.

Method	Cases
<b>Image method</b>	$G(x, s) = \frac{1}{2\pi} \left\{ \ln x-s  + \lim_{N \rightarrow \infty} \left[ \left( \sum_{m=1}^N -\ln x-s_{4i-3}  - \ln x-s_{4i-2}  + \ln x-s_{4i-1}  + \ln x-s_{4i}  \right) + c_1(N)\ln x-s_{c1}  + c_2(N)\ln x-s_{c2}  + e(N) \right] \right\}$
	$G(x, s) = \frac{1}{2\pi} \left\{ \ln x-s  + \lim_{N \rightarrow \infty} \left[ \left( \sum_{i=1}^N \ln x-s_{2i-1}  + \ln x-s_{2i}  \right) + c_1(N)\ln x-s_{c1}  + c_2\ln x-s_{c2}  + d_1(N) \left[ \ln x-s_{d1}  + \sum_{j=1}^M \ln x-s_j^1  \right] + d_2(N) \left[ \ln x-s_{d2}  + \sum_{j=1}^M \ln x-s_j^2  \right] \right] \right\}$
<b>Bipolar coordinates</b>	$x_1 - d = \frac{a^2}{x_2 - d}, \quad x_2 = \frac{b^2}{x_1 + d}, \quad x_1 = \frac{a^2}{x_2}, \quad x_2 = d - x_1, \text{ when } b = a$ $c = \frac{(x_2 - x_1)}{2} = \frac{\sqrt{a^4 - 2a^2b^2 + b^4 - 2a^2d^2 - 2b^2d^2 + d^4}}{2d} \Rightarrow \frac{\sqrt{d^2 - 4a^2}}{2} (a = b)$ $c = \frac{\sqrt{a^4 - 2a^2b^2 + b^4 - 2a^2d^2 - 2b^2d^2 + d^4}}{2d}$



**Fig. 4.** Final images and the focuses of the bipolar coordinate (a) an eccentric annulus, (b) a half plane with a circular hole and (c) an infinite plane with two circular holes.

The problem sketch of an eccentric annulus is shown in Fig. 2. The location of image source and bipolar coordinates are shown in Fig. 4(a). The source point is located at  $s=(0,0.75)$ . The centers of two holes are set at  $(0,0)$  and  $(-0.4,0)$ , and the radii are 0.4 and 1.0 for the inner and outer boundaries, respectively. Following the success of annulus case for the iterative images, we now extend to the eccentric case. In a similar way of finding the successive images for matching the inner and outer boundary conditions [18], the solution can be superimposed using Eq.(20). Finally, we can find that the final frozen image points and the focuses of the bipolar coordinates are the same. After collocating some points to

match the boundary conditions, all the unknown coefficients can be determined. The results are compared well with the analytical solution using the bipolar coordinates. The contour plots using the present method of Eq. (20), the bipolar coordinates of Eq. (13) and the null-field BIEM [19] are shown in Fig. 5.

**Case 2.** A half plane containing a circular hole

Fig. 4(b) depicts Green's function for the half plane with a hole and the homogeneous Dirichlet boundary condition. The source point is located at  $s=(3,0)$ . The center and radius of the hole is  $(0,0)$  and  $a=1$ , respectively. The  $d/2=1.25$  is the distance from the

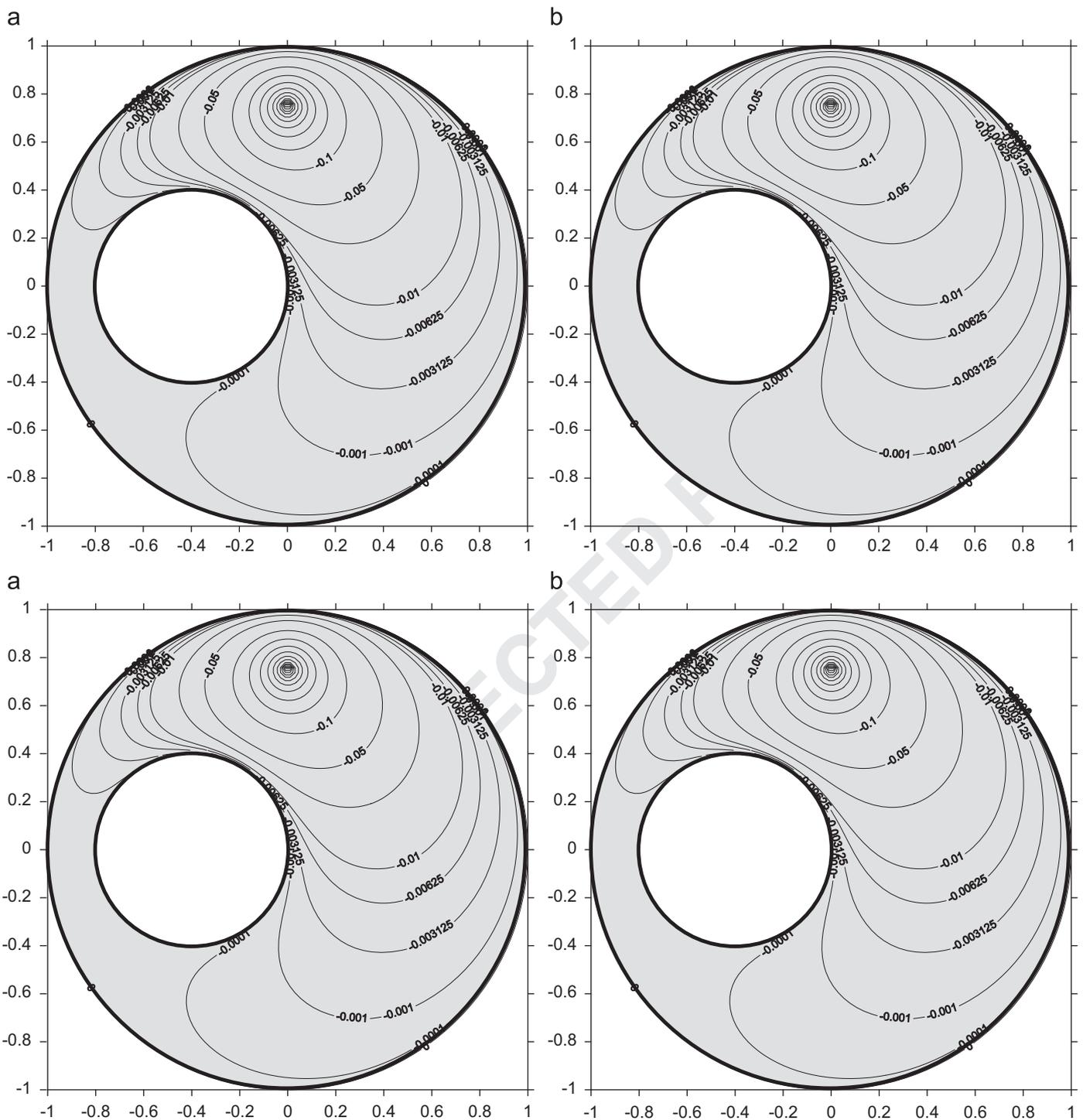


Fig. 5. Green's function using (a) an analytical solution using bipolar coordinate, (b) image solution, (c) solution using superposition technique and (d) solution using Green's third identity in the null-field BIEM [18]

center to the ground line. Similarly, the analytical and semi-analytical solutions are obtained using the bipolar coordinates and the image method, respectively. The results agree well with those of the null-field BIEM [19] in Fig. 6.

### Case 3. An infinite plane containing two circular holes

Following the successful experiences of the eccentric annulus case for the iterative images, we now extend to the infinite plane containing two circular holes as shown in Table 1. The

problem sketch of the infinite plane containing two circular holes is shown in Fig. 4(c). The source point is located at  $s=(3.85,0)$ . The centers of two holes are set at  $(0,0)$  and  $(2.1,0)$ , and their radii are 0.4 and 1, respectively. In a similar way of finding the image sources for matching boundary conditions [18], an image solution is derived.

We also found that the final frozen image points approach to the foci of the bipolar coordinates. Based on the image solution for an infinite plane containing a circular hole subject to the

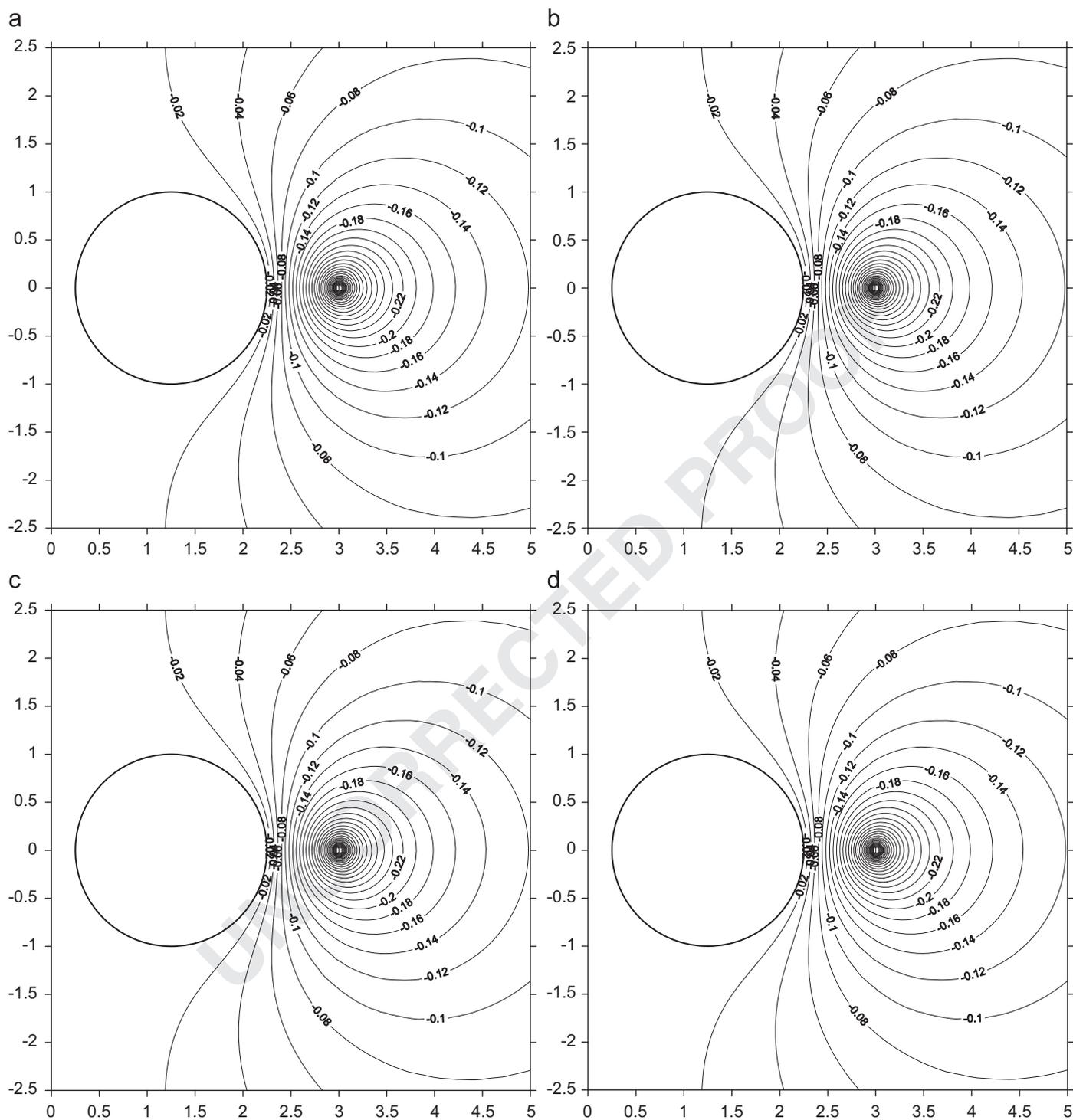


Fig. 6. Green's function using (a) an analytical solution using bipolar coordinate, (b) image solution, (c) solution using superposition technique and the null-field BIEM [18] and (d) solution using Green's third identity in the null-field BIEM.

homogeneous Neumann BC, an extra source at the center of hole is required. This motivates us to put sources at two centers of the holes to obtain acceptable results. Therefore, Eq. (19) is extended to

$$G(x,s) = \frac{1}{2\pi} \left\{ \ln|x-s| + \lim_{N \rightarrow \infty} \left[ \sum_{i=1}^N \ln|x-s_{2i-1}| + \ln|x-s_{2i}| \right] + c_1(N) \ln|x-s_{c1}| + c_2(N) \ln|x-s_{c2}| \right\}$$

$$\left. \begin{aligned} &+ d_1(N) \left[ \ln|x-s_{d1}| + \sum_{j=1}^M \ln|x-s_j^1| \right] \\ &+ d_2(N) \left[ \ln|x-s_{d2}| + \sum_{j=1}^M \ln|x-s_j^2| \right] \end{aligned} \right\}, \quad (21)$$

where two extra sources  $s_{d1}$  and  $s_{d2}$  are located at the two centers of holes,  $s_j^1$  and  $s_j^2$  are the successive images due to  $s_{d1}$  and  $s_{d2}$ , respectively. The results agree well with those of the null-field BIEM [19] and the conventional MFS in Fig. 7. It is

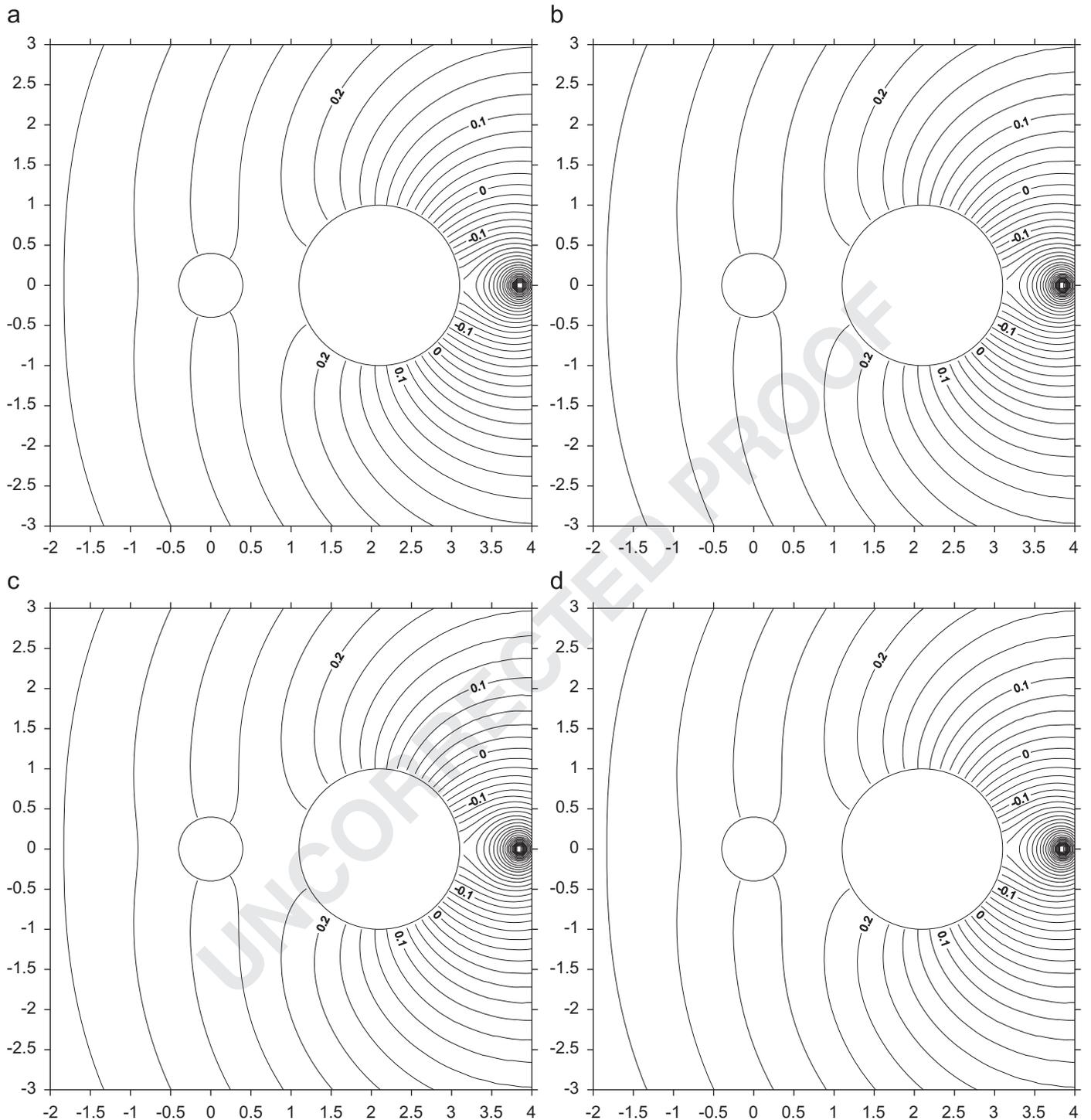


Fig. 7. Green's function using (a) an analytical solution using bipolar coordinate, (b) image solution, (c) solution using superposition technique and the null-field BIEM [18] and (d) solution using Green's third identity in the null-field BIEM.

interesting to find that the final images also freeze at focuses in the bipolar coordinates. The results are summarized in Table 1.

For arbitrary boundaries, it is not easy to have such a neat semi-analytical solution. Furthermore, the simple case is extension to a straight boundary and can be found in textbooks. Our approach can also be easily extended to solve inclusion problems by taking free body. One part is a boundary value problem without a source, the other is also a BVP with a source. Additionally, matching the boundary conditions on the interface is required.

## 5. Conclusions

In this paper, Green's functions were derived using the image method. It is found that final image points terminate at the two focuses of the bipolar coordinates for all the three cases, an eccentric annulus, a half plane containing a circular hole and an infinite plane containing two circular holes. The optimal source distribution in the MFS is dependent on the given geometry and the source location. An image method can guide as to search for an optimal source location of the MFS and can determine the

strengths of sources except the two frozen images. Three examples were demonstrated to find all the image sources for constructing Green's function. The dimension of the influence matrix in the linear algebraic equation is at most four by four in all the examples. Agreement is made after comparing with other solutions.

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